

New Exact Traveling Wave Solutions for the Zakharov Equations and the Coupled Klein-Gordon-Zakharov Equations

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Abstract: The modified simplest equation method is an efficient method for obtaining exact solutions of some nonlinear partial differential equations. In the present work, the modified simplest equation method is used to construct exact solutions of the differential equations. In the present work, the modified simplest equation method is used to construct exact solutions of the Zakharov equations and the coupled Klein-Gordon-Zakharov equations. It also shown that the proposed method is effective and general.

Keywords: Modified simplest equation method, Zakharov equations, coupled Klein-Gordon-Zakharov equations.

1. Introduction

Nonlinear systems of partial differential equations (PDEs) play an important role in various areas of modern physics and engineering such as fluid mechanics, plasma physics, optical fibers and quantum mechanics. In theoretical investigation of the dynamics of strong Langmuir turbulence in plasma physics, various Zakharov equations take an important role [1, 2]. In this paper, we consider the following Zakharov equations

$$\begin{cases} n_{tt} - c_s^2 n_{xx} = \beta(|E|^2)_{xx}, \\ iE_t + \alpha E_{xx} = \delta nE, \end{cases}$$

and the coupled Klein-Gordon-Zakharov equations:

$$\begin{cases} u_{tt} - c_0^2 \nabla^2 u + f_0^2 u + \delta u v = 0, \\ v_{tt} - c_0^2 \nabla^2 v - \beta \nabla^2 |u|^2 = 0. \end{cases}$$

Zakharove equations and coupled Klein-Gordon-Zakharov equations are model systems to describe nonlinear interactions in plasma, where E or u denotes the electric field (more precisely it denotes its slowly varying envelope), n or v denotes the ion density fluctuation.

More recently, some exact solutions for the Zakharov and the coupled Klein-Gordon-Zakharov equations are obtained by using different methods [3-18]. In this work, we apply the modified simplest equation method to the Zakharov and the coupled Klein-Gordon-Zakharov equations. The modified simplest equation method is one ofthe most powerful and direct methods for constructing solutions of nonlinear partial differential equations.

In this paper, we have proposed the modified simplest equation method, and have presented applications for this method to nonlinear partial differential equations. The rest of the paper is organized as follows. In section 2, we describe the modified simplest equation method for finding traveling wave solutions of nonlinear evolution equations, and give the main steps of the method. In the subsequent sections, we will apply the method to find exact traveling wave solutions of the Zakharov and the coupled Klein-Gordon-Zakharov equations. In the last section, some conclusions are presented.

2. Description of the modified simplest equation method

In this section, we briefly review the modified simplest equation method [19-23]. That is based on the assumption that the exact solutions can be expressed by a polynomial in $\frac{F'}{F}$ such that $F = F(\xi)$ is an unknown linear ordinary equation to be determined later. This method consists of the following steps:

Step 1. Consider a general form of nonlinear partial differential equation (PDE)

$$P(u, u_x, u_t, u_{xx}, u_{tx}, \dots) = 0$$
(1)

Assume that the solution is given by $u(x,t) = U(\xi)$ where $\xi = x + ct$. Hence, we use the following changes:

we use the following changes:

$$\frac{\partial}{\partial t}(0) = c \frac{\partial}{\partial \xi}(0),$$

$$\frac{\partial}{\partial x}(0) = \frac{\partial}{\partial \xi}(0),$$

$$\frac{\partial^2}{\partial x^2}(0) = \frac{\partial^2}{\partial \xi^2}(0)$$
(2)

and so on for other derivatives. Using (2) changes the PDE (1) to an ODE

$$Q(U, U'U'', ...) = 0.$$
 (3)

where $U = U(\xi)$ is an unknown function, Q is a polynomial in the variable U and its derivatives.

Step 2.We suppose that Eq. (3) has the following formal solution:

$$U(\xi) = \sum_{i=0}^{N} A_i \left(\frac{F'}{F}\right)^i, \tag{4}$$

where A_i are arbitrary constants to be determined such that $A_N \neq 0$ while $F(\xi)$ is an unknown function to be determined later.

Step 3. We determine the positive integer N in (4) by balancing the highest order derivatives and the nonlinear terms in Eq.(3).

Step 4. We substitute (4) into (3), we calculate all the necessary derivatives U, U', U'', \dots and then we account the



function $F(\xi)$. As a result of this substitution, we get a polynomial of $\frac{F'(\xi)}{\pi C}$ and its derivatives. polynomial, we equate all the coefficients to zero. This operation yields a system of equations which can be solved to find A_i and $F(\xi)$. Consequently, we can get the exact solution of Eq.(1).

3. Application the modified simplest equation method

In this section, we study the Zakharov and the coupled Klein-Gordon-Zakharov equation using the modified simplest equation method.

The modified simplest equation method to the 3.1 Zakharov equations

Consider the Zakharov equations

$$\begin{cases} n_{tt} - c_s^2 n_{xx} = \beta(|E|^2)_{xx}, \\ iE_t + \alpha E_{xx} = \delta nE, \end{cases}$$
 (5)

$$iE_t + \alpha E_{xx} = \delta nE, \tag{6}$$

which is one of the classical models governing dynamics of nonlinear waves and describing interactions between highand low-frequency waves, where n is the perturbed number density of the ion (in the low-frequency response), E is slow variation amplitude of the electric field intensity, c_s is the thermal traity transportation velocity of the electron-ion, and $\alpha \neq 0, \beta \neq 0, \delta \neq 0$ and c_s are constants.

Since E(x, t) in (6) is a complex function we assume that

$$E(x,t) = U(\xi) \exp[i(kx - \omega t + \xi_0)],$$

$$n(x,t) = n(\xi), \quad \xi = x - c_g t + \xi_1,$$
(7)

where $U(\xi)$ is a real-valued function c_g , k and w are constants to be determined later, and ξ_0 and ξ_1 are constants. Substituting (7) into Eqs. (5), (6), we have

$$(c_g^2 - c_s^2) n'' = \beta (U^2)''$$

$$\alpha U'' + i (2\alpha k - c_g) U' + (\omega - \alpha k^2) U - \delta n U = 0.$$
(8)

Integrating (8) twice with respect to ξ and taking integration constants to zero yields

$$n = \frac{\beta}{c_g^2 - c_s^2} U^2$$
, $c_g^2 - c_s^2 \neq 0$. (10)

We assume that when $c_g^2 > c_s^2$ (supersonic speed), n is positive; and when $c_q^2 < c_s^2$ (subsonic speed), n is negative

Substituting (10) into (9),

$$\alpha U'' + i \left(2\alpha k - c_g \right) U' + (\omega - \alpha k^2) U - \frac{\beta \delta}{c_g^2 - c_s^2} U^3 = 0.$$
(11)

In view of (9) we assume that

$$c_q = 2\alpha k$$
, $\gamma = \omega - \alpha k^2$.

Then Eq. (11) becomes the nonlinear ODE

$$U''(\xi) + k_3 U^3(\xi) + k_1 U(\xi) = 0, \tag{12}$$

where

$$k_3 = \frac{\delta \beta}{\alpha (4\alpha^2 k^2 - c_s^2)}, \qquad k_1 = \frac{\gamma}{\alpha}.$$

By balancing the highest order derivative term U'' with the nonlinear term U^3 in (12), we obtain N=1 in (4). So

we assume that Eq. (12) has solution in the form

$$U(\xi) = A_0 + A_1 \left(\frac{F'}{F}\right), \quad A_1 \neq 0.$$
 (13)

Using (13), we obtain

$$U^{3} = A_{0}^{3} + 3A_{0}^{2}A_{1}\left(\frac{F'}{F}\right) + 3A_{0}A_{1}^{2}\left(\frac{F'}{F}\right)^{2} + A_{1}^{3}\left(\frac{F'}{F}\right)^{3},$$
(14)

$$U'' = A_1 \left(\frac{F'''}{F} - 3 \frac{F'F''}{F^2} + 2 \left(\frac{F'}{F} \right)^3 \right). \tag{15}$$

Substituting (13) to (15) into Eq. (12) and setting the coefficients of F^{j} (j = 0, -1, -2) to zero, we obtain

$$k_1 A_0 + k_3 A_0^3 = 0, (16)$$

$$A_1 F''' + k_1 A_1 F' + 3k_3 A_0^2 A_1 F' = 0, (17)$$

$$-3A_1F'F'' + 3k_3A_0A_1^2F'^2 = 0, (18)$$

$$-3A_1F'F'' + 3k_3A_0A_1^2F'^2 = 0,$$

$$2A_1F'^3 + k_3A_1^3F'^3 = 0.$$
(18)

Eqs. (16) and (19) directly imply following solutions:

$$A_0 = \pm \sqrt{-\frac{k_1}{k_3}}, \quad A_1 = \pm \sqrt{-\frac{2}{k_3}}, \quad k_1 > 0, \quad k_3 < 0.$$

Thus, Eqs. (17) and (18) become

$$F''' - 2k_1 F' = 0, (20)$$

$$-F'' + \sqrt{2k_1}F' = 0. (21)$$

By substituting Eq. (21) into Eq. (20) we get

$$-\sqrt{2k_1}F'' + F''' = 0. (22)$$

The general solution of Eq. (22) is

$$F(\xi) = a_0 + a_1 \xi + a_2 \exp(\sqrt{2k_1}\xi),$$

where a_i (i = 0,1,2) are arbitrary constants. Thus, we have

$$U(\xi) = \pm \sqrt{-\frac{k_1}{k_3}} \pm \sqrt{-\frac{2}{k_3}} \left(\frac{a_1 + \sqrt{2k_1}a_2 \exp(\sqrt{2k_1}\xi)}{a_0 + a_1\xi + a_2 \exp(\sqrt{2k_1}\xi)} \right),$$

and

$$n(\xi) = \frac{\beta}{c_g^2 - c_s^2} \left(\pm \sqrt{-\frac{k_1}{k_3}} \right)$$
$$\pm \sqrt{-\frac{2}{k_3} \left(\frac{a_1 + \sqrt{2k_1}a_2 \exp(\sqrt{2k_1}\xi)}{a_0 + a_1\xi + a_2 \exp(\sqrt{2k_1}\xi)} \right)^2}.$$

Now, the exact solution of Eqs.(5) and (6) have the form

$$\begin{split} &E(x,t) \\ &= \pm \sqrt{-\frac{k_1}{k_3}} \\ &\pm \sqrt{-\frac{2}{k_3}} \left(\frac{a_1 + \sqrt{2k_1}a_2 \exp\left(\sqrt{2k_1}(x - c_g t + \xi_1)\right)}{a_0 + a_1(x - c_g t + \xi_1) + a_2 \exp\left(\sqrt{2k_1}(x - c_g t + \xi_1)\right)} \right) \end{split}$$

$$\times \exp(i(kx - \omega t + \xi_0)).$$
 (23)

and



$$n(x,t) = \frac{\beta}{c_g^2 - c_s^2}$$

$$\times \left(\pm \sqrt{-\frac{k_1}{k_3}} \pm \sqrt{-\frac{2}{k_3}}\right)$$

$$\left(\frac{a_1 + \sqrt{2k_1}a_2 \exp(\sqrt{2k_1}(x - c_g t + \xi_1))}{a_0 + a_1(x - c_g t + \xi_1) + a_2 \exp(\sqrt{2k_1}(x - c_g t + \xi_1))}\right)$$
(24)

If $a_1 = 0$ and $a_0 = a_2 = 1$, we have

$$E(x,t) = \pm \sqrt{-\frac{k_1}{k_3}} \left(2 + \tanh \sqrt{\frac{k_1}{2}} \left(x - c_g t + \xi_1 \right) \right)$$
$$\times \exp(i(kx - \omega t + \xi_0))$$

and

$$n(x,t) = \frac{\beta}{c_g^2 - c_s^2} \left[\sqrt{-\frac{k_1}{k_3}} (2 + \tanh \sqrt{\frac{k_1}{2}} (x - c_g t + \xi_1)) \right]^2$$

Example. Solve the Zakharov equations by using the odified simplest equation method

$$\begin{cases} n_{tt} - n_{xx} = (|E|^2)_{xx}, \\ iE_t + E_{xx} = nE, \end{cases}$$

 $\begin{cases} n_{tt}-n_{xx}=(|E|^2)_{xx},\\ iE_t+E_{xx}=nE,\\ \text{Substituting } c_s^2=1,\ \beta=1,\ \alpha=1\ \text{and}\ \delta=1\ \text{in (23) and} \end{cases}$ (24) gives

$$\begin{split} &E(x,t) = \pm \sqrt{(\omega - k^2)(1 - 4k^2)} \pm \sqrt{2(1 - 4k^2)} \\ &\times \left(\frac{a_1 + \sqrt{2(\omega - k^2)}a_2 \exp\left(\sqrt{2(\omega - k^2)}(x - c_g t + \xi_1)\right)}{a_0 + a_1(x - c_g t + \xi_1) + a_2 \exp\left(\sqrt{2(\omega - k^2)}(x - c_g t + \xi_1)\right)} \right) \\ &\times exp\big(i(kx - \omega t + \xi_0)\big). \end{split}$$

and

$$n(x,t) = \frac{1}{4k^2 - 1} \left[\pm \sqrt{(\omega - k^2)(1 - 4k^2)} \pm \sqrt{2(1 - 4k^2)} \right]$$

$$\times \left(\frac{a_1+\sqrt{2(\omega-k^2)}a_2\exp\left(\sqrt{2(\omega-k^2)}\left(x-c_gt+\xi_1\right)\right)}{a_0+a_1\big(x-c_gt+\xi_1\big)+a_2\exp\left(\sqrt{2(\omega-k^2)}\left(x-c_gt+\xi_1\right)\right)}\right)]^2.$$

3.2 The modified simplest equation method to the coupled Klein-Gordon-Zakharov equations

In this subsection, we consider the coupled nonlinear Klein-Gordon-Zakharov equations

$$\begin{cases} u_{tt} - c_0^2 \nabla^2 u + f_0^2 u + \delta u v = 0 \\ v_{tt} - c_0^2 \nabla^2 v - \beta c_0^2 \nabla^2 |u|^2 = 0. \end{cases}$$
 (25)

We seek its following wave packet solution

$$u(x, y, z, t) = U(\xi) \exp(i(kx + ly + nz - \Omega t)),$$

$$v(x, y, z, t) = V(\xi), \quad \xi = px + qy + rz - \omega t \tag{26}$$

where both $U(\xi)$ and $V(\xi)$ are real-valued functions k, l, n, Ω, p, q, r and ω are constants to be determined later. Substituting Eq. (26) into Eq. (25) yields

$$(w^{2}-c_{0}^{2}P^{2})U''(\xi) + 2i(\omega\Omega-c_{0}^{2}KP)U'(\xi)$$

$$-(w^{2}-c_{0}^{2}K^{2}-f_{0}^{2})U(\xi) + \delta V(\xi)U(\xi) = 0,$$

$$(w^{2}-c_{0}^{2}P^{2})V''(\xi) - \beta P^{2}(U^{2}(\xi))'' = 0.$$
(27)

$$K = (k, l, n), \quad K^2 = k^2 + l^2 + n^2, \quad P = (p, q, r), \quad P^2$$

= $p^2 + q^2 + r^2, \quad K.P = kp + lq + nr$

In view of (27) we assume that

$$\omega\Omega = c_0^2 K. P, \tag{28}$$

then (27) is reduced to

$$(w^{2} - c_{0}^{2}P^{2})U''(\xi) - (w^{2} - c_{0}^{2}K^{2} - f_{0}^{2})U(\xi) + \delta V(\xi)U(\xi) = 0,$$
 (29)

$$(w^2 - c_0^2 P^2) V''(\xi) - \beta P^2 (U^2(\xi))'' = 0, \tag{30}$$

Integrating (30) once respect to
$$\xi$$
 we get $(w^2 - c_0^2 P^2)V'(\xi) - \beta P^2(U^2(\xi))' = c_1,$ (31)

where c_1 is integration constant. Because we find the special form of exact solutions for simplicity purpose, we take $c_1 =$ 0 and integrating this formula once again, we have

$$V(\xi) = \frac{c_2}{(w^2 - c_0^2 P^2)} + \frac{\beta P^2}{(w^2 - c_0^2 P^2)} U^2(\xi), \tag{32}$$

where c_2 is integration constant. Substituting (32) into (29)

$$\begin{split} &(w^2 - c_0^2 P^2)^2 U''(\xi) + \left[(w^2 - c_0^2 P^2) \left((w^2 - c_0^2 K^2 - f_0^2) \right) + \\ &\delta c_2 \right] U(\xi) + \delta \beta P^2 U^3(\xi) = 0. \end{split} \tag{33} \\ &Eq. \ (33) \ \text{can be expressed as} \end{split}$$

$$U''(\xi) + k_3 U^3(\xi) + k_1 U(\xi) = 0$$
 (34)

where

$$k_1 = \frac{\left[(w^2 - c_0^2 P^2) \left((w^2 - c_0^2 K^2 - f_0^2) \right) + \delta c_2 \right]}{(w^2 - c_0^2 P^2)^2},$$

$$k_3 = \frac{\delta \beta P^2}{(w^2 - c_0^2 P^2)^2}$$

By balancing the highest order derivative term U'' with the nonlinear term U^3 in (34), we obtain N=1 in (4). So we assume that Eq. (34) has solution in the form

$$U(\xi) = A_0 + A_1 \left(\frac{F'}{F}\right), \quad A_1 \neq 0.$$
 (35)

Using (35), we obtain

$$U^{3} = A_{0}^{3} + 3A_{0}^{2}A_{1}\left(\frac{F'}{F}\right) + 3A_{0}A_{1}^{2}\left(\frac{F'}{F}\right)^{2} + A_{1}^{3}\left(\frac{F'}{F}\right)^{3}, (36)$$

$$U'' = A_1 \left(\frac{F'''}{F} - 3 \frac{F'F''}{F^2} + 2 \left(\frac{F'}{F} \right)^3 \right). \tag{37}$$

Substituting (35) to (37) into Eq. (34) and setting the coefficients of F^{j} (j = 0, -1, -2) to zero, we obtain

$$k_1 A_0 + k_3 A_0^3 = 0, (38)$$

$$k_1 A_0 + k_3 A_0^3 = 0,$$
 (38)
 $A_1 F''' + k_1 A_1 F' + 3k_3 A_0^2 A_1 F' = 0,$ (39)



$$-3A_1F'F'' + 3k_3A_0A_1^2F'^2 = 0,$$

$$2A_1F'^3 + k_3A_1^3F'^3 = 0.$$
(40)

Eqs. (38) and (41) directly imply following solutions:

$$A_0 = \pm \sqrt{-\frac{k_1}{k_3}}, \quad A_1 = \pm \sqrt{-\frac{2}{k_3}}, \quad k_1 > 0, \quad k_3 < 0.$$

Thus, Eqs. (39) and (40) become

$$F''' - 2k_1F' = 0, (42)$$

$$-F'' + \sqrt{2k_1}F' = 0. (43)$$

By substituting Eq. (43) into Eq. (42) we get $-\sqrt{2k_1}F^{\prime\prime}+F^{\prime\prime\prime}=0. \tag{44}$

The general solution of Eq. (44) is

$$F(\xi) = a_0 + a_1 \xi + a_2 \exp(\sqrt{2k_1}\xi)$$

where a_i (i = 0,1,2) are arbitrary constants.

Thus, we have

$$U(\xi) = \pm \sqrt{-\frac{k_1}{k_3}} \pm \sqrt{-\frac{2}{k_3}} \left(\frac{a_1 + \sqrt{2k_1}a_2 \exp(\sqrt{2k_1}\xi)}{a_0 + a_1\xi + a_2 \exp(\sqrt{2k_1}\xi)} \right),$$

and

$$\begin{split} V(\xi) &= \frac{c_2}{(w^2 - c_0^2 P^2)} + \frac{\beta P^2}{(w^2 - c_0^2 P^2)} \\ &\times (\pm \sqrt{-\frac{k_1}{k_3}} \pm \sqrt{-\frac{2}{k_3}} \left(\frac{a_1 + \sqrt{2k_1} a_2 \exp(\sqrt{2k_1} \xi)}{a_0 + a_1 \xi + a_2 \exp(\sqrt{2k_1} \xi)} \right))^2. \end{split}$$

Now, the exact solution of Eq. (25) has the form n(x,y,z,t)

$$\begin{split} &u(x,y,z,t)\\ &=\pm\sqrt{-\frac{k_{1}}{k_{3}}}\pm\sqrt{-\frac{2}{k_{3}}}\\ &\times\left(\frac{a_{1}+\sqrt{2k_{1}}a_{2}\exp\left(\sqrt{2k_{1}}(px+qy+rz-\omega t)\right)}{a_{0}+a_{1}(px+qy+rz-\omega t)+a_{2}\exp\left(\sqrt{2k_{1}}(px+qy+rz-\omega t)\right)}\right) \end{split}$$

$$\times \exp(i(kx + ly + nz - \Omega t)). \tag{45}$$

and

$$v(x, y, z, t) = \frac{c_2}{(w^2 - c_0^2 P^2)} + \frac{\beta P^2}{(w^2 - c_0^2 P^2)} \left[\pm \sqrt{-\frac{k_1}{k_3}} \pm \sqrt{-\frac{2}{k_3}} \right] \times \left(\frac{a_1 + \sqrt{2k_1} a_2 \exp\left(\sqrt{2k_1} (px + qy + rz - \omega t)\right)}{a_0 + a_1 (px + qy + rz - \omega t) + a_2 \exp\left(\sqrt{2k_1} (px + qy + rz - \omega t)\right)} \right)^2$$

$$(46)$$

If $a_1 = 0$ and $a_0 = a_2 = 1$, we have

$$u(x, y, z, t) = \pm \sqrt{-\frac{k_1}{k_3}} (2 + \tanh \sqrt{\frac{k_1}{2}} (px + qy + rz - \omega t))$$
$$\times \exp(i(kx + ly + nz - \Omega t))$$

and

$$\begin{split} v(x,y,z,t) &= \frac{c_2}{(w^2 - c_0^2 P^2)} + \frac{\beta P^2}{(w^2 - c_0^2 P^2)} \\ &\times \left[\pm \sqrt{-\frac{k_1}{k_3}} (2 + \tanh\sqrt{\frac{k_1}{2}} (px + qy + rz - \omega t)) \right]^2 \end{split}$$

Example. Solve the Klein-Gordon-Zakharov equations by using the modified simplest equation method

$$\begin{cases} u_{tt} - u_{xx} + u + uv = 0, \\ v_{tt} - v_{xx} - (|u|^2)_{xx} = 0. \end{cases}$$

Substituting $c_0^2 = 1, f_0^2 = 1, \beta = 1, \delta = 1$ and $c_2 = 0$ in (45) and (46) gives

$$\begin{split} u(x,y,z,t) &= \pm \sqrt{\frac{(\Omega^2 - k^2 - 1)}{(\omega^2 - p^2)}} \pm \sqrt{-\frac{2(\omega^2 - p^2)^2}{p^2}} \\ \times \left(\frac{a_1 + \sqrt{2(\frac{-\Omega^2 + k^2 + 1}{\omega^2 - p^2})} a_2 \exp\left(\sqrt{2(\frac{-\Omega^2 + k^2 + 1}{\omega^2 - p^2})} (px + qy + rz - \omega t)\right)}{a_0 + a_1(px + qy + rz - \omega t) + a_2 \exp\left(\sqrt{2(\frac{-\Omega^2 + k^2 + 1}{\omega^2 - p^2})} (px + qy + rz - \omega t)\right)} \right) \\ \times \exp\left(i(kx + ly + nz - \Omega t)\right). \end{split}$$

 $v(x, y, z, t) = \frac{p^2}{(w^2 - p^2)} \left[\pm \sqrt{\frac{(\Omega^2 - k^2 - 1)}{(\omega^2 - p^2)}} \pm \sqrt{-\frac{2(\omega^2 - p^2)^2}{p^2}} \right]$ $\left\{ \frac{a_1 + \sqrt{2(\frac{-\Omega^2 + k^2 + 1}{\omega^2 - p^2})} a_2 \exp\left(\sqrt{2(\frac{-\Omega^2 + k^2 + 1}{\omega^2 - p^2})} (px + q + rz - \omega t)\right)}{a_0 + a_1(px + qy + rz - \omega t) + a_2 \exp\left(\sqrt{2(\frac{-\Omega^2 + k^2 + 1}{\omega^2 - p^2})} (px + qy + rz - \omega t)\right)} \right]^2$

4. Conclusion

In this work, we obtained exact solutions of the Zakharov and the coupled Klein-Gordon-Zakharov equations by using the modified simplest equation method. The results show that this method is efficient.

References

- [1] S. G. Thornhill, D. Haar, "Langmuir turbulence and modulational instability", *Phys. Rep*, vol. 43, no. 2, pp. 43-99, 1978.
- [2] R. O. Dendy, Plasma Dynamics, *Oxford University Press*, Oxford, 1990.
- [3] J. L. Zhang, M. L. Wang, D. M. Chen, Z. D. Fang, "The periodic wave solutions for two nonlinear evolution equations", Commun. Theor. Phys, vol. 40, no. 2, pp. 129-132, 2003.
- [4] Naranmandula, K. X. Wang, "New spiky and explosive solitary wave solutions for further modified Zakharov-Kuznetsov equation", *Phys. Lett. A*, vol. 336, no. 2–3, pp. 112-116, 2005.
- [5] C. H. Zhao, Z. M. Sheng, "Explicit traveling wave solutions for Zakharov equation", *Acta Phys. Sinica*, vol. 53, no. 6, pp. 1629-1634, 2004.
- [6] D. J. Huang, H. Q. Zhang, "Extended hyperbolic function method and new exact solitary wave solutions of Zakharov equations", Acta Phys. Sinica, vol. 53, no. 8, pp. 2434-2438, 2004.



- [7] Y. Chen, B. Li, "New exact traveling wave solutions for generalized Zakharov-Kuzentsov equations using general projective Riccati equation method', *Commun. Theor. Phys.*, vol. 41, no. 1, pp. 1-6, 2004.
- [8] B. Li, Y. Chen, H. Q. Zhang, "Exact traveling wave solutions for a generalized Zakharov-Kuznetsov equation", *Appl. Math. Comput*, vol. 46, no. 2-3, pp. 653-666, 2003.
- [9] Y. Chen, B. Li, H. Q. Zhang, "Bäcklund transformation and exact solutions for a new generalized Zakharov-Kuzentsov equation with nonlinear terms of any order", Commun. Theor. Phys, vol. 39, no. 2, pp. 135-140, 2003.
- [10] A. Biswas, E. Zerrad, J. Gwanmesia and R. Khouri, "1-soliton solution of the generalized Zakharov equation in plasmas by He,s variational principle", *Applied Mathematics and Computation*, vol. 215, no. 12, pp. 4462-4466, 2010.
- [11] M. S. Ismail and A. Biswas, "1-soliton solution of the Klein-Gordon-Zakharov equation with power law nonlinearity", *Applied Mathematics and Computation*, vol. 217, no. 8, pp. 4186-4196, 2010.
- [12] P. Suarez and A. Biswas, "Exact 1-soliton solution of the Zakharov equation in plasmas with power law nonlinearity", *Applied Mathematics and Computation*, vol. 217, no. 17, pp. 7372-7375, 2011.
- [13] G. Ebadi, E. V. Krishnan and A. Biswas, "Solitons and cnoidal waves of the Klein-Gordon Zakharov equation in plasma", *Pramana*, vol. 79, no. 2, pp. 185-198, 2012.
- [14] M. Song, B. Ahmed and A. Biswas, "Topological soliton solution and bifurcation analysis of the Klein-Gordon-Zakharov equation in (1+1)-dimensions with power law nonlinearity", *Journal of Applied Mathematics*, 2013, 972416, 2013.
- [15] R. Morris, A. H. Kara and A. Biswas, "Soliton solution and conservation laws of the Zakharov equation in plasmas with

- power law nonlinearity'', Nonlinear Analysis: Modelling and Control, vol. 18, no. 2, pp. 153-159, 2013.
- [16] B. S. Ahmed, E. Zerrad and A. Biswas, "Kinks and domain walls of the Zakharov equation in plasmas", *Proceedings of* the Romanian Academy, Series A, vol. 14, no. 4, pp. 281-286, 2013.
- [17] M. Song, B. Ahmed, E. Zerrad and A. Biswas, "Domain wall and bifurcation analysis of the Klein-Gordon-Zakharov equation in (2+1)-dimensions with power law nonlinearity", *Chaos*, vol. 23, no. 3, 033115, 2013.
- [18] A. H. Bhrawy, M. A. Abdelkawy, S. Kumar, S. Johnson and A. Biswas, "Solitons and other solutions to quantum Zakharov-Kuznetsov equation in quantum magneto plasmas", *Indian Journal of Physics*, vol. 87, no. 5, pp. 455-463, 2013.
- [19] A. J. M. Jawad, M. D. Petkovic, A. Biswas, "Modified simple equation method for nonlinear evolution equations", *Appl Math Comput*, vol. 217, no. 2, pp. 869-877, 2010.
- [20] M. E. Z. Elsayed, "A note on the modified simple equation method applied to Sharma-Tasso-Olver equation", Appl Math Comput, vol. 218, no. 7, pp. 3962-3964, 2011.
- [21] K. V. Nikolay, I. D. Zlatinka, H. Kantz, "Modified method simplest equation and application to nonlinear PDEs", *Appl Math Comput*, vol. 216, pp. 2587-2595, 2010.
- [22] K. V. Nikolay, "Modified method simplest equation powerful tool for obtaining exact and approximate traveling-wave solutions of nonlinear PDEs", Commun Nonlinear Sci Numer Simulat, vol. 16, pp. 1179-1185, 2011.
- [23] K. V. Nikolay, I. D. Zlatinka, "Application of the method of simplest equation for obtaining exact traveling-wave solutions for two classes of model PDEs from ecoloy and population dynamics", Commun Nonlinear Sci Numer Simulat, vol. 15, pp. 2836-2845, 2010.